

The slope at point B,

$$\begin{aligned}\tan(\theta - \delta\theta) &\simeq \theta - \delta\theta \\ &= \frac{\partial y}{\partial x} - \frac{\delta \ddot{y}}{\partial x^2} \delta x\end{aligned}$$

The upward component of the tension at B = $T \tan(\theta - \delta\theta)$

$$\begin{aligned}&= T(\theta - \delta\theta) \\ &= T \left[\frac{\partial y}{\partial x} - \frac{\delta \ddot{y}}{\partial x^2} \delta x \right] \\ &\rightarrow \textcircled{2}\end{aligned}$$

The resultant downward tension [from eqⁿ ① and ②]

$$\begin{aligned}F &= T \frac{\partial y}{\partial x} - T \left[\frac{\partial y}{\partial x} - \frac{\delta \ddot{y}}{\partial x^2} \delta x \right] \\ &= T \frac{\delta \ddot{y}}{\partial x^2} \delta x \rightarrow \textcircled{3}\end{aligned}$$

Let m be the mass of the string per unit length.

\therefore Mass of the element AB = $m \delta x$

Acceleration of the element along y axis —

$$= \frac{\delta \ddot{y}}{\delta t^2}$$

\therefore Force acting on the element, $F = m \delta x \frac{\delta \ddot{y}}{\delta t^2} \rightarrow \textcircled{4}$

$$[F = ma]$$

From eqⁿ ③ and ④,

$$m \delta x \frac{\delta \ddot{y}}{\delta t^2} = T \frac{\delta \ddot{y}}{\delta x^2} \delta x$$

So, the volume enclosed within A, to B,

$$= (dx + d\frac{V}{V}) \cdot 1$$

$$= dx + d\frac{V}{V}$$

Change in volume undergone by the original volume \rightarrow

$$= (dx + d\frac{V}{V}) - dx$$

$$= d\frac{V}{V}$$

\therefore Change in volume per unit volume,

$$= \frac{d\frac{V}{V}}{dx} = \text{volume strain}$$

Now, bulk modulus, ~~strain~~

$$E = \frac{\text{Stress}}{\text{Volume strain}}$$

$$\Rightarrow \text{Stress} = E \times \text{volume strain}$$

$$= E \times \frac{d\frac{V}{V}}{dx}$$

$$\Rightarrow \text{Pressure} = E \times \frac{d\frac{V}{V}}{dx}$$

If P be the pressure on A and $P + \frac{dP}{dx} \cdot dx$ be the pressure on B then the change in pressure =

$$P + \frac{dP}{dx} dx - P = \frac{dP}{dx} dx$$

$$\Rightarrow \frac{\delta^2 y}{\delta t^2} = \frac{T}{m} \frac{\delta^2 y}{\delta x^2} \rightarrow (5)$$

This is similar to the differential equation of wave motion,

$$\frac{d^2 y}{dt^2} = v^2 \frac{d^2 y}{dx^2} \rightarrow (6)$$

Comparing eqⁿ (5) and (6),

$$v^2 = \frac{T}{m}$$

$$\therefore v = \sqrt{\frac{T}{m}} \rightarrow (7)$$

We know,

$$v = \lambda \nu, \text{ where } \nu = \text{frequency}$$

If l be the length of the string for p segment then the length of each segment = $\frac{l}{p}$

$$\therefore \frac{l}{p} = \frac{\lambda}{2} \text{ (for one segment)}$$

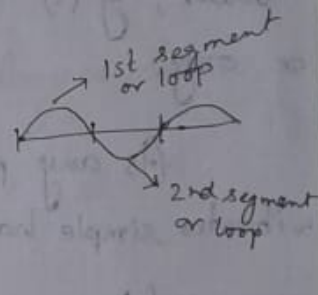
$$\therefore \lambda = \frac{2l}{p}$$

$$\therefore v = \lambda \nu \left(\frac{2l}{p} \right) \rightarrow (8)$$

From (7) and (8),

$$\sqrt{\frac{T}{m}} = \nu \left(\frac{2l}{p} \right)$$

$$\therefore \nu = \frac{p}{2l} \sqrt{\frac{T}{m}}$$



$$= \frac{d}{dx} \left(E \frac{d\xi}{dx} \right) dx$$

$$= E \frac{d^2\xi}{dx^2} dx$$

= Force per unit area.

$$= \rho dx \frac{d^2\xi}{dt^2}$$

$$P = \frac{F}{A}$$

$$= \frac{ma}{A}$$

$$= \frac{v \rho a}{A}$$

$$= dx \rho a$$

$$= \rho dx \frac{d^2\xi}{dt^2}$$

Hence, $E \frac{d^2\xi}{dx^2} dx = \rho dx \frac{d^2\xi}{dt^2}$

$$\therefore \frac{d^2\xi}{dt^2} = \frac{E}{\rho} \frac{d^2\xi}{dx^2} \rightarrow \textcircled{1}$$

We know,

the expression of differential equation of wave

motion, $\frac{d^2\xi}{dt^2} = v^2 \frac{d^2\xi}{dx^2} \rightarrow \textcircled{2}$

Comparing equation $\textcircled{1}$ and $\textcircled{2}$,

$$v^2 = \frac{E}{\rho}$$

$$\Rightarrow v = \sqrt{\frac{E}{\rho}}$$

Thus, the velocity of the sound in gaseous medium

$$v = \sqrt{\frac{E}{\rho}}$$

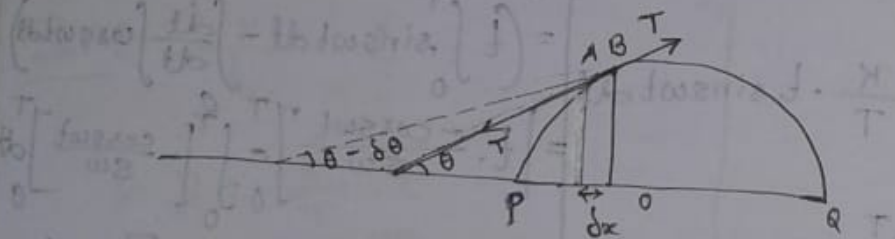
where, E = elasticity of the medium.

ρ = density of the medium.

Equation of transverse vibration of a stretched string.

or

Velocity of transverse vibration of a stretched string.



Consider a string PQ stretched by a tension T along x direction. Let the string is plucked at the centre O and let it free. So the string vibrates along y direction.

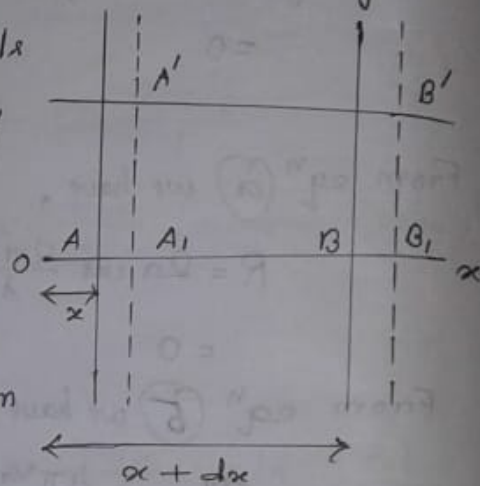
Consider a small element AB of length Δx . The tangents drawn at A and B makes angles θ and $\theta - \Delta\theta$

Sound waves

Velocity of Longitudinal waves in Gases (Fluid)

Imagine a longitudinal plane sound wave consisting of condensation and rarefaction travels through a gas medium of density ρ to the right with velocity c .

Let ox be the direction of advance of the wave and let AB be a very thin layer of the medium of unit cross section.



Let x be the distance of A from the origin at any instant and $x + dx$ be the distance of B from the origin.

Hence the thickness of the slab $AB = dx$

the volume enclosed with the slab $AB = dx$.

Let ξ be the displacement of A to A_1 after time dt and $\xi + d\xi$ be the displacement of B to B_1 .

Now, the distance between the phase,

$$A_1 \text{ to } B_1 = (x + dx + \xi + d\xi) - (x + \xi) = dx + d\xi$$

with the x axis.

The tension at A is resolved into two rectangular components. The downward component of tension at $A = T \sin \theta$.

If θ is small,

$$\therefore \sin \theta \approx \theta$$

Again, $\theta \approx \tan \theta$

Now, $\tan \theta = \frac{dy}{dx}$ at A (i.e. slope at A)

$$\therefore \theta = \frac{dy}{dx}$$

\therefore The downward component of tension at $A = T \tan \theta$

$$= T \frac{dy}{dx} \quad \rightarrow \textcircled{1}$$

The rate of change of slope w.r.t the length of the element = $\frac{d}{dx} \left(\frac{dy}{dx} \right)$

$$= \frac{d^2y}{dx^2}$$

Thus, the change in slope for a distance, δx

$$\delta \theta = \frac{d^2y}{dx^2} \delta x$$